

broken links
energy ϵ per broken link
 g internal states for each half

a) $g=1$ mean number \bar{n} of broken links

$$k_B T \gg \epsilon$$

$$F=0$$

$$\bar{n}d \ll L$$

partition function $Z = \sum_{n=0}^N e^{-\beta n \epsilon} \approx \sum_{n=0}^{\infty} e^{-\beta n \epsilon}$

$$Z = \sum e^{-\beta \epsilon}$$

$$= \frac{1}{1 - e^{-\beta \epsilon}}$$

$$\frac{\partial Z}{\partial \beta} = \sum -\epsilon e^{-\beta \epsilon}$$

$$\log Z = -\log(1 - e^{-\beta \epsilon})$$

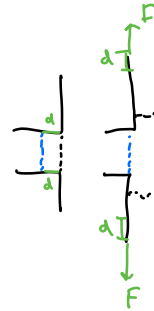
$$\langle E \rangle = -\frac{1}{Z} \frac{\partial Z}{\partial \beta} = -\frac{\partial \log Z}{\partial \beta}$$

$$-\frac{\partial}{\partial \beta} \log Z = \frac{1}{1 - e^{-\beta \epsilon}} \cdot -e^{-\beta \epsilon} \cdot -\epsilon = \frac{\epsilon e^{-\beta \epsilon}}{1 - e^{-\beta \epsilon}}$$

$$\rightarrow \bar{n} = \frac{\langle E \rangle}{\epsilon} = \frac{e^{-\beta \epsilon}}{1 - e^{-\beta \epsilon}} = \frac{1}{e^{\beta \epsilon} - 1} \approx \frac{1}{1 + \frac{\epsilon}{k_B T} - 1} = \frac{k_B T}{\epsilon}$$

$$\boxed{\bar{n} = \frac{k_B T}{\epsilon}}$$

$k_B T \gg \epsilon$



When a force is applied, the energy per link becomes $\epsilon' = \epsilon - 2Fd$

$$\bar{n} \rightarrow \frac{k_B T}{\epsilon'} = \boxed{\frac{k_B T}{\epsilon - 2Fd}}$$

b) $g > 1$

partition function

$$Z = \sum_{n=0}^N (g^2)^n e^{-\beta n \epsilon'} \approx \sum_{n=0}^{\infty} (g^2)^n e^{-\beta n \epsilon'}$$

internal states

$$= \sum_{n=0}^{\infty} g^{2n} e^{-\beta n (\epsilon - 2Fd)}$$

$$Z = \frac{1}{1 - g^2 e^{-\beta(\epsilon - 2Fd)}}$$

free energy

$$G = -k_B T \log Z = k_B T \log [1 - g^2 e^{-\beta(\epsilon - 2Fd)}]$$

At $T = T_c$, $Z \rightarrow \infty$, $G \rightarrow -\infty$ since there will be an "infinite" number of broken links

$$1 - g^2 e^{-\beta_c(\epsilon - 2Fd)} = 0$$

$$e^{-\beta_c(\epsilon - 2Fd)} = \frac{1}{g^2}$$

$$-\beta_c(\epsilon - 2Fd) = -\log g^2$$

$$\frac{\epsilon - 2Fd}{k_B T_c} = 2 \log g \rightarrow T_c = \frac{\epsilon - 2Fd}{2k_B \log g}$$

$$c) Z = \frac{1}{1 - g^2 e^{-\beta(\epsilon - 2Fd)}}$$

$$\bar{n} = \frac{\langle E \rangle}{\epsilon'} = -\frac{1}{\epsilon'} \frac{\partial \log Z}{\partial \beta} = \frac{1}{\epsilon'} \frac{\partial}{\partial \beta} \log [1 - g^2 e^{-\beta(\epsilon - 2Fd)}]$$

$$= \frac{1}{\epsilon'} \frac{1}{1 - g^2 e^{-\beta(\epsilon - 2Fd)}} \cdot -g^2 e^{-\beta(\epsilon - 2Fd)} \cdot -(\epsilon - 2Fd)$$

$$\bar{h} = \frac{\varepsilon - 2Fd}{\underbrace{\varepsilon'}_{\downarrow}} \frac{g^2 e^{-\beta(\varepsilon - 2Fd)}}{1 - g^2 e^{-\beta(\varepsilon - 2Fd)}}$$

$$\bar{h} = \frac{1}{g^2 e^{\beta(\varepsilon - 2Fd)} - 1}$$